

233B HW #1

1. Synchrotron radiation loss

According to the formula (1), the synchrotron radiation loss scales as E^4/R^2 , once $v \approx c$. However, in order to compensate for the falling cross section $\sigma \propto 1/E^2$, we need to increase $\mathcal{L} \propto E^2$, which requires $N_{\pm} \propto E$ according to Eq. (2) if f_c is kept constant. Therefore, the overall synchrotron loss scales as E^5/R^2 to maintain the same event rate for annihilation processes. Of course, one could have chosen to increase the repetition rate $f_c \propto E^2$, but that would require the number of bunches to increase as E^2 , and the synchrotron radiation loss as E^6/R^2 . Therefore, increasing the number of electrons in the beam is more beneficial (if it is possible to maintain the beam stable in an entirely different manner). Using $R = 27$ km for $E = 100$ GeV ($\sqrt{s} = 200$ GeV for LEP-II), we scale it up to $R = 27 \text{ km} * (2 \text{ TeV} / 200 \text{ GeV})^{5/2} = 8500 \text{ km}$. This is clearly impractical.

There is currently only one known way to overcome this problem in designing higher energy $e^+ e^-$ colliders. You make it completely linear, with two beams colliding head-on from two LINACs (LINEar ACcelerators). This is the concept behind the ILC (International Linear Collider).

$$\mathbf{N} [27 * 10^{5/2}]$$
$$8538.15$$

2. α_s from J/ψ decay

According to the PDG particle listing, the branching fractions of J/ψ decays are

$$\begin{aligned} \text{inclusive hadronic:} & \quad 87.7 \pm 0.5 \% \\ & \quad \text{among which virtual photon contribution is } 13.50 \pm 0.30 \% \\ \mu^+ \mu^-: & \quad 5.93 \pm 0.06 \% \end{aligned}$$

Separating out the $g g g$ contribution (non-virtual-photon piece), we find $74.2 \pm 0.6 \%$ (assumed uncorrelated Gaussian error, which may not be too conservative.) Taking the ratio to $\mu^+ \mu^-$, we find

$\frac{\Gamma(J/\psi \rightarrow g g g)}{\Gamma(J/\psi \rightarrow \mu^+ \mu^-)} = 12.5 \pm 0.16$. Equating it to the lowest order prediction, we find $\alpha_s = 0.205 \pm 0.001$. The error due to the higher order corrections is actually much bigger than this error bar which we ignore. (Actually at this energy, α is more like $1/134$, which changes the result very little.) This result is not crazy (but a little low), compared to Fig. 9.2 in Ian Hinchliffe's review on QCD in PDG, <http://pdg.lbl.gov/2007/reviews/qcdrpp.pdf>

As a cross check, we can see if the hadronic width through the virtual photon agrees with the lowest order formula. We find $\frac{\Gamma(J/\psi \rightarrow q \bar{q})}{\Gamma(J/\psi \rightarrow \mu^+ \mu^-)} = 2.28$, which should be compared to $N_c \left(\left(\frac{2}{3}\right)^2 + \left(\frac{-1}{3}\right)^2 + \left(\frac{-1}{3}\right)^2 \right) = 2$, which is in a *reasonably* good agreement given that we ignore the higher order QCD corrections. We see that there are 10-20% errors because we ignored the higher order corrections.

$$\mathbf{87.7 - 13.5}$$
$$74.2$$

$$\text{Sqrt}[0.5^2 + 0.3^2]$$

$$0.583095$$

$$74.2 / 5.93$$

$$12.5126$$

$$\% * \text{Sqrt}\left[\left(\frac{0.6}{74.2}\right)^2 + \left(\frac{0.06}{5.93}\right)^2\right]$$

$$0.162068$$

$$\left(\frac{16\pi \left(\frac{2}{3}\right)^2 \left(\frac{1}{137}\right)^2}{\frac{160}{81} (\pi^2 - 9)} 12.51\right)^{1/3}$$

$$0.205423$$

$$\frac{13.50}{5.93}$$

$$2.27656$$

3. J/ψ line shape

The cross section into $g g g$ final state is given by the Breit-Wigner formula because there is no interference, while that into $q \bar{q}$ final state interferes between the resonant and virtual photon contributions. The Breit-Wigner contribution is

$$\sigma_{\text{BW}} = \frac{2J+1}{(2S_1+1)(2S_2+1)} \frac{4\pi}{k^2} \frac{m^2 \Gamma^2}{(s-m^2)^2 + m^2 \Gamma^2} B_{\text{in}} B_{\text{out}} = \frac{12\pi}{s} \frac{m^2 \Gamma^2}{(s-m^2)^2 + m^2 \Gamma^2} B_{\text{in}} B_{\text{out}}$$

The case of the $\mu^+ \mu^-$ final state has the interference and the cross section is

$$\sigma = \frac{4\pi\alpha^2}{3s} \left| 1 + \frac{\xi^2 s}{s-m^2 + im\Gamma} \right|^2$$

To fix ξ , we drop the non-resonant piece and require it gives the Breit-Wigner formula,

$$\frac{4\pi\alpha^2}{3s} \frac{\xi^4 s^2}{(s-m^2)^2 + m^2 \Gamma^2} = \frac{12\pi}{s} \frac{m^2 \Gamma^2}{(s-m^2)^2 + m^2 \Gamma^2} B_{\text{in}} B_{\text{out}}$$

and hence (using $s \approx m^2$),

$$\alpha^2 \xi^4 m^2 = 9 \Gamma^2 B_{ee} B_{\mu\mu}$$

or

$$\alpha \xi^2 m = 3 \Gamma B_{ll}$$

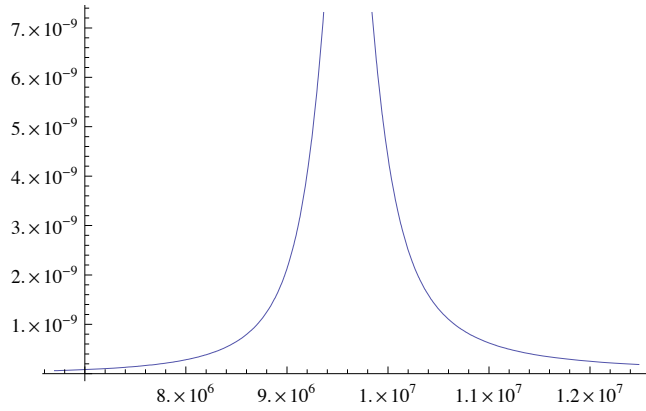
Having fixed ξ , the cross section to the $q \bar{q}$ final state is

$$\sigma = \frac{4\pi\alpha^2}{3s} \left| 1 + \frac{3m\Gamma B_{ll}}{\alpha} \frac{1}{s-m^2 + im\Gamma} \right|^2 \frac{B_{qq}}{B_{ll}}$$

$$\text{Plot}\left[\frac{86.8}{s 10^6} \text{Abs}\left[1 + \frac{3 m \Gamma B_{11}}{\alpha} \frac{1}{s - m^2 + I m \Gamma}\right]^2 \frac{B_{qq}}{B_{11}} / .\right.$$

$$\left.\{m \rightarrow 3096.916, \Gamma \rightarrow 93.4, \alpha \rightarrow \frac{1}{137}, B_{11} \rightarrow 0.0594, B_{qq} \rightarrow 0.1350\},\right.$$

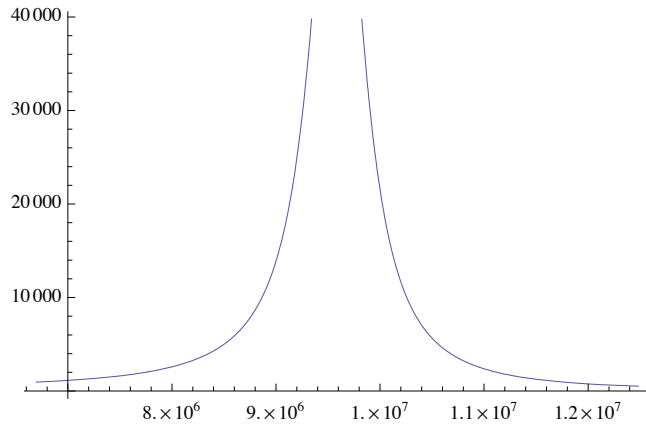
$$\left.\{s, 3096.916^2 - 10 \times 3096.916 \times 93.4, 3096.916^2 + 10 \times 3096.916 \times 93.4\}\right]$$



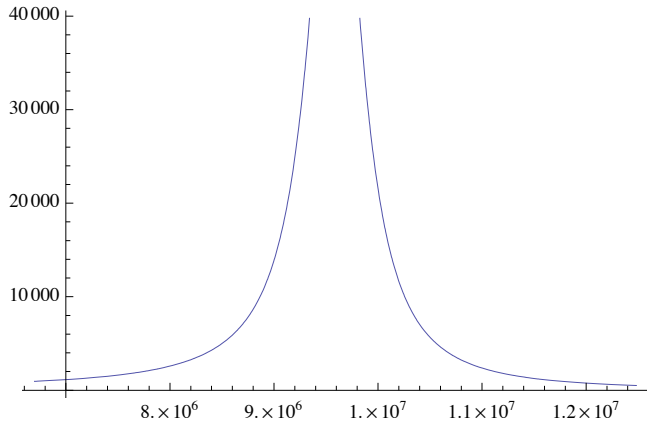
$$\text{Plot}\left[0.389 10^6 \frac{12 \pi}{s 10^{-6}} \text{Abs}\left[\frac{m \Gamma}{s - m^2 + I m \Gamma}\right]^2 B_{ggg} B_{11} / .\right.$$

$$\left.\{m \rightarrow 3096.916, \Gamma \rightarrow 93.4, \alpha \rightarrow \frac{1}{137}, B_{11} \rightarrow 0.0594, B_{qq} \rightarrow 0.1350, B_{ggg} \rightarrow 0.877 - 0.1350\},\right.$$

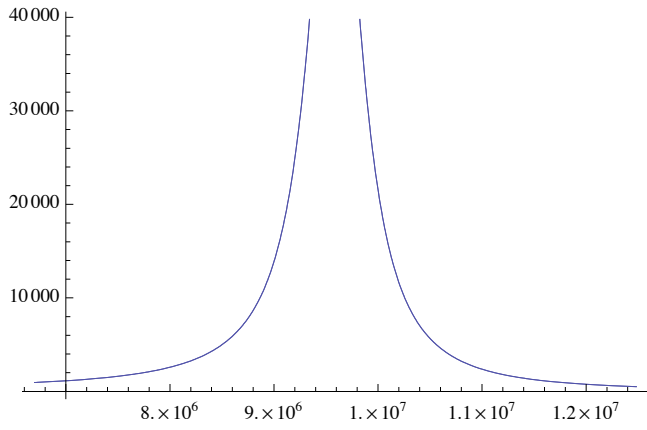
$$\left.\{s, 3096.916^2 - 10 \times 3096.916 \times 93.4, 3096.916^2 + 10 \times 3096.916 \times 93.4\}\right]$$



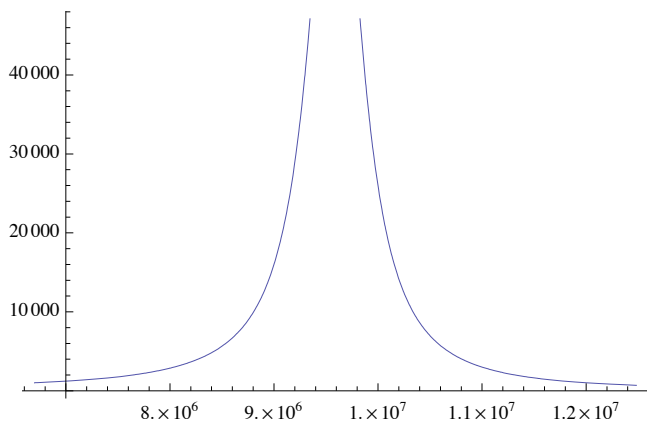
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Plot[ $\frac{86.8}{s 10^{-6}} \text{Abs}\left[\frac{3 m \Gamma B_{11}}{\alpha} \frac{1}{s - m^2 + I m \Gamma}\right]^2 \frac{B_{ggg}}{B_{11}}$  /.
  {m -> 3096.916, \Gamma -> 93.4, \alpha -> \frac{1}{137}, B_{11} -> 0.0594, B_{qq} -> 0.1350, B_{ggg} -> 0.877 - 0.1350},
  {s, 3096.916^2 - 10 \times 3096.916 \times 93.4, 3096.916^2 + 10 \times 3096.916 \times 93.4}]
```



Show[%, %%]



$$\text{Plot}\left[\frac{86.8}{s 10^{-6}} \text{Abs}\left[1 + \frac{3 m \Gamma B_{11}}{\alpha} \frac{1}{s - m^2 + I m \Gamma}\right]^2 \frac{B_{qq}}{B_{11}} + \frac{86.8}{s 10^{-6}} \text{Abs}\left[\frac{3 m \Gamma B_{11}}{\alpha} \frac{1}{s - m^2 + I m \Gamma}\right]^2 \frac{B_{ggg}}{B_{11}}\right] /. \left\{m \rightarrow 3096.916, \Gamma \rightarrow 93.4, \alpha \rightarrow \frac{1}{137}, B_{11} \rightarrow 0.0594, B_{qq} \rightarrow 0.1350, B_{ggg} \rightarrow 0.877 - 0.1350\right\}, \left\{s, 3096.916^2 - 10 \times 3096.916 \times 93.4, 3096.916^2 + 10 \times 3096.916 \times 93.4\right\}]$$



$$\text{Plot}\left[\text{NIntegrate}\left[\frac{1}{\sqrt{2 \pi} \sigma} E^{-\left(\sqrt{s-t}\right)^2 / 2 / \sigma^2} \frac{86.8}{s 10^{-6}} \text{Abs}\left[1 + \frac{3 m \Gamma B_{11}}{\alpha} \frac{1}{s - m^2 + I m \Gamma}\right]^2 \frac{B_{qq}}{B_{11}} + \frac{86.8}{s 10^{-6}} \text{Abs}\left[\frac{3 m \Gamma B_{11}}{\alpha} \frac{1}{s - m^2 + I m \Gamma}\right]^2 \frac{B_{ggg}}{B_{11}}\right] /. \left\{m \rightarrow 3096.916, \Gamma \rightarrow 93.4, \alpha \rightarrow \frac{1}{137}, B_{11} \rightarrow 0.0594, B_{qq} \rightarrow 0.1350, B_{ggg} \rightarrow 0.877 - 0.1350\right\} /. \left\{\sigma \rightarrow 0.90\right\}, \left\{s, (t - 10)^2, (t + 10)^2\right\}], \left\{t, 1546 * 2, 1551 * 2\right\}\right]$$

