

D)

Py 229C

22/11/2005

Make up class

Dark Matter: So far!

We concluded dark matter might be Wimps,
 Cold Thermal relics,
 looked at annihilation cross section
 New physics at TeV scale
 Detection experiments etc.

Alternative ideas about dark matter:

- Non thermal relics

① Thermal relic

↓ decays

Something lighter

Abundance $\Omega_{DM} \propto \frac{1}{\sigma_{ann} \text{ of DM}}$

$$\hookrightarrow \propto \frac{1}{\sigma_{ann} \text{ of parent}} \cdot \frac{M_{DM}}{M_{parent}}$$

Ex:

SUSY

 $\tilde{\chi}^0$ is the parent particle $\hookrightarrow \gamma \tilde{G}_{\text{gravitino}}$ linkage between $\Omega_{DM} + \sigma_{ann}^{DM}$ is broken.

2)

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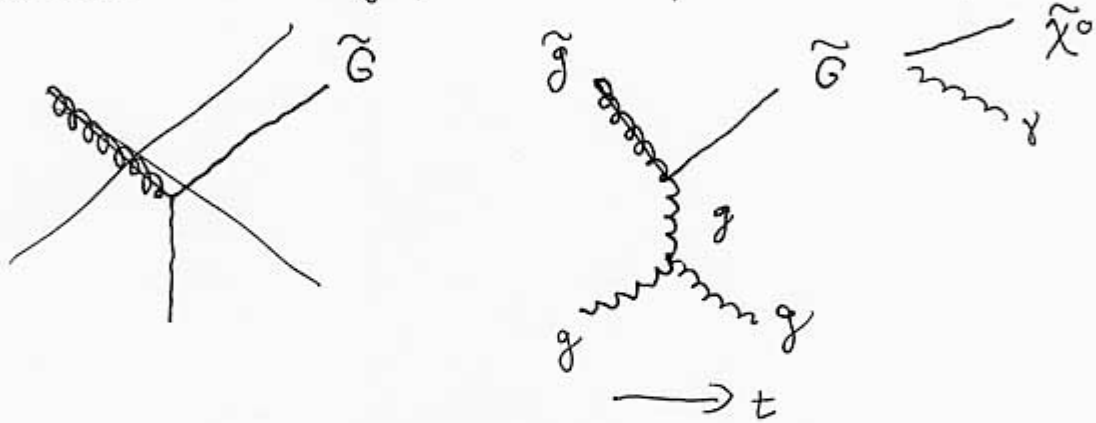
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② Thermally produced in non-equilibrium



thermal bath, assume no \tilde{G}

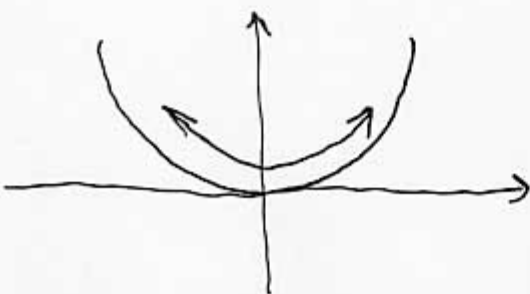
Calculate the # of gravitino produced in the bath.



$$\Omega_{\tilde{G}} \approx \frac{T_{RH}}{10^{16} \text{ GeV}} \frac{M_{\tilde{G}}}{100 \text{ GeV}}$$

reheating temperature
= "highest T of the universe ever"
(See later inflation)

③ Scalar Condensate



harmonic oscillator

$$K_e \sim \frac{1}{2} \dot{x}^2$$

$$P_e \sim \frac{1}{2} x^2$$

Virial Theorem:

$$\langle K \rangle = \langle V \rangle$$

$\langle \rangle =$ time average

(3)

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22/11/2005

relativistic field

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right]$$

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & R^2(t) & 0 & 0 \\ 0 & 0 & R^2(r) & 0 \\ 0 & 0 & 0 & R^2(\theta) \end{pmatrix} \quad \text{assuming flat universe } k=0$$

If ϕ has mass m , if its self interaction can be ignored if its interaction with anything else can be ignored

$$\Rightarrow V(\phi) = \frac{1}{2} m^2 \phi^2$$

\Rightarrow Spin 0 particle of mass m , real $\phi \leftrightarrow$ particle = anti particle

Energy momentum tensor

$$\begin{aligned} T^{\mu\nu} &= \partial^\mu \phi \partial^\nu \phi - g^{\mu\nu} \mathcal{L} \\ &= \partial^\mu \phi \partial^\nu \phi - g^{\mu\nu} \left(\frac{1}{2} g^{\rho\sigma} \partial_\rho \phi \partial_\sigma \phi - \frac{1}{2} m^2 \phi^2 \right) \end{aligned}$$

$\phi(\vec{x}, t) = \phi(t)$ spatially homogenous configuration

$$\mathcal{L} = \frac{1}{2} g^{\rho\sigma} \partial_\rho \phi \partial_\sigma \phi - \frac{1}{2} m^2 \phi^2 = \frac{1}{2} \dot{\phi}^2 - \frac{1}{2} m^2 \phi^2$$

$$T = \begin{pmatrix} \frac{1}{2} \dot{\phi}^2 + \frac{1}{2} m^2 \phi^2 & 0 & 0 & 0 \\ 0 & -\frac{1}{R^2} (\frac{1}{2} \dot{\phi}^2 - m^2 \phi^2) & 0 & 0 \\ 0 & 0 & " & 0 \\ 0 & 0 & 0 & " \end{pmatrix}$$

4)

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$$T^{\mu\nu} = \begin{pmatrix} \rho & & & \\ & p & & \\ & & p & \\ & & & p \end{pmatrix} \text{ for fluid}$$

$$\text{so } \rho = \frac{1}{2} \dot{\phi}^2 + \frac{1}{2} m^2 \phi^2$$

$$p = \frac{1}{2} \dot{\phi}^2 - m^2 \phi^2 \quad \langle p \rangle = 0 \text{ by virial Theorem}$$

$$w = p/\rho = 0 \quad (\text{same as non-relativistic gas of massive particles})$$

$$\rho \propto R^{-3(1+w)} = R^{-3}$$

GR derivation

Field equation of ϕ

$$\frac{1}{\sqrt{-g}} \partial_\nu \sqrt{-g} g^{\mu\nu} \partial_\mu \phi + V'(\phi) = 0$$

$$g^{\mu\nu} = \begin{pmatrix} 1 & & & \\ & -k^2 & & \\ & & -R^2 & \\ & & & -R^2 \end{pmatrix}$$

$$\frac{1}{R^3} \partial_t R^3(t) g^{00} \partial_t \phi + m^2 \phi = 0$$

$$\ddot{\phi} + 3 \frac{\dot{R}}{R} \dot{\phi} + m^2 \phi = 0$$

$$\ddot{\phi} + 3H(t) \dot{\phi} + m^2 \phi = 0$$

assume $H(t)$ is a constant when we solve this.
This is justified if $m \gg H$

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Solve using Fourier transform:

$$\phi = \int d\omega \tilde{\phi}(\omega) e^{-i\omega t}$$

$$= -\omega^2 \tilde{\phi} + 3H(-i\omega) \tilde{\phi}(\omega) + m^2 \tilde{\phi}(\omega) = 0$$

$$-\omega^2 - (3H\omega + m^2) = 0$$

$$\text{So } \omega = \frac{1}{2} \left(-3iH \pm \sqrt{-H^2 + 4m^2} \right)$$

$$= \pm m - \frac{3}{2} iH + O(H^2/m)$$

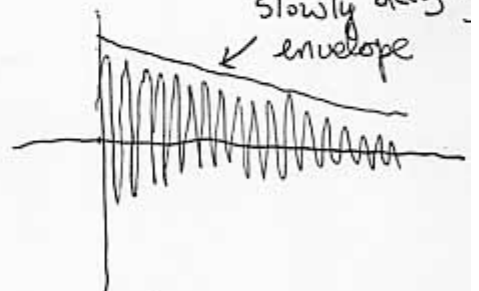
$$\text{So } \phi(t) = \tilde{\phi}(\omega) e^{i\omega t}$$

$$= \tilde{\phi}(\omega) e^{\mp i\omega t} e^{-\frac{3}{2} H t}$$

$$\phi(t) = \cos(mt + \varphi_0) e^{-\frac{3}{2} H t}$$

as long as $m \gg H$

slowly decaying envelope

but if H is varying with t we can

$$\text{write: } e^{-\frac{3}{2} H t} = e^{-\frac{3}{2} \int H(t) dt}$$

$$= e^{-\frac{3}{2} \int \dot{R}/R dt} = e^{-\frac{3}{2} \log R} = R^{-3/2}$$

$$\text{So } \phi(t) = \tilde{\phi} \cos(mt + \varphi_0) R^{-3/2}$$

So harmonic oscillator of a scalar field behaves ~~exactly~~ as non-relativistic matter

\Rightarrow Dark Matter ???

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Classically oscillating field = coherent state of particles

$a^\dagger(\vec{p}=0)$ creation operator

$$|f\rangle \equiv e^{f a^\dagger} |0\rangle$$

$$\langle f | \phi(\vec{x}, t) | f \rangle = f \cos mt$$

How can this be dark matter?

assume $\phi(t_{\text{early}}) \neq 0$

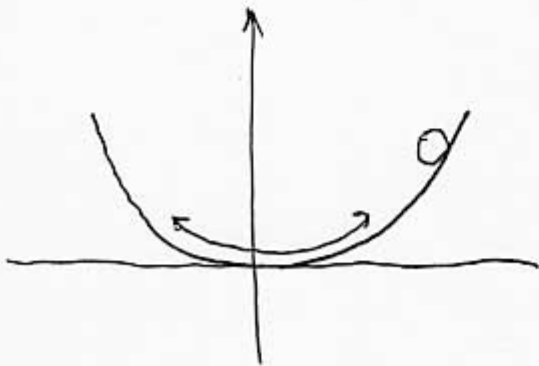
$$\omega^2 + 3iH\omega - m^2 = 0$$

$H \gg m$

$$\omega = \frac{1}{2} \left(-3iH \pm \sqrt{-9H^2 + 4m^2} \right)$$

$$= \cancel{3iH} + 0 \quad \hookrightarrow -3iH + o\left(\frac{m^2}{H}\right)$$

$$\omega = \begin{cases} -3iH & \pm \left(\frac{m^2}{H}\right) e^{-3Ht} \\ 0 & e^0 \end{cases}$$



When $H \gg m$

ϕ doesn't move

$H \sim m$

ϕ starts to move

$H \ll m$

ϕ : matter

Most popular candidate along these lines; axion

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Motivation for axion

Strong interactions described by Quantum Chromodynamics

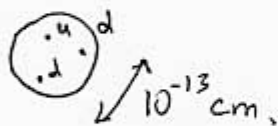
$$\mathcal{L} = -\frac{1}{4g^2} G_{\mu\nu}^a G^{\mu\nu a} + \theta \frac{1}{64\pi^2} G_{\mu\nu}^a G_{\rho\sigma}^a \epsilon^{\mu\nu\rho\sigma}$$

↓
 $\vec{E} \cdot \vec{B}$

Second term violates P, I

neutron electric dipole moment = 0

$$H = \vec{S} \cdot \vec{E} \quad d_n \leq 10^{-26} \text{ e cm.}$$



limit on neutron EDM $\rightarrow \theta < 10^{-10}$

could have been $O(1)$
but instead is extremely
small!

Must be a physical
mechanism that sets
 $\theta = 0$?

Introduce the axion:

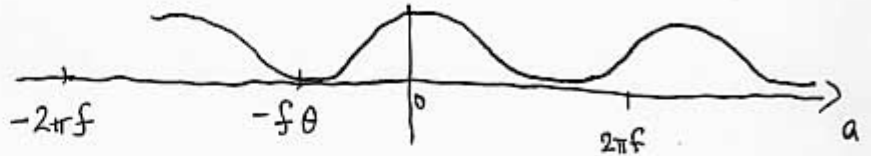
Instead of thinking of ~~theta~~ as a parameter
think of θ as a dynamical field:
real scalar field a :

(8)

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$$\left(\theta + \frac{a}{f}\right) \frac{1}{64\pi^2} G_{\mu\nu}^a G^{\mu\nu a} \epsilon^{\mu\nu\rho\sigma}$$



at minimum a cancels f
 \rightarrow no \cancel{p} , \cancel{X} , dn

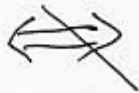
f axion decay constant [E]

$$m_a = 0.62 \times 10^{-3} \text{ eV} \times \frac{10^{10} \text{ GeV}}{f}$$

When f is low α_f GG is strong, proton/neutron star
 would emit too many axions ($f \lesssim 10 \text{ GeV}$)

SN
1987A

cools too fast



neutrino burst: from SN 1987A
 observed \geq a few seconds.

~~Limit~~ on axion density

$$\Omega_a h^2 \approx \underbrace{1.9 \times 3^{\pm 1}}_{O(1) \text{ for } f \sim 10^{12} \text{ GeV}} \left(\frac{1 \text{ eV}}{m_a} \right)^{1.175} \underbrace{\rho_i^2 F(\rho_i)}_{\text{initial condition}}$$

(9)

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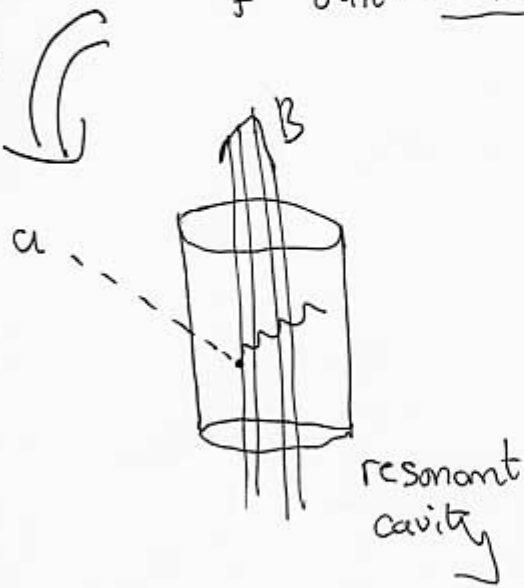
Axion could couple to E.M. field:

$$g/f \frac{1}{64\pi} G_{\mu\nu}^a G_{\rho\sigma}^a \epsilon^{\mu\nu\rho\sigma}$$

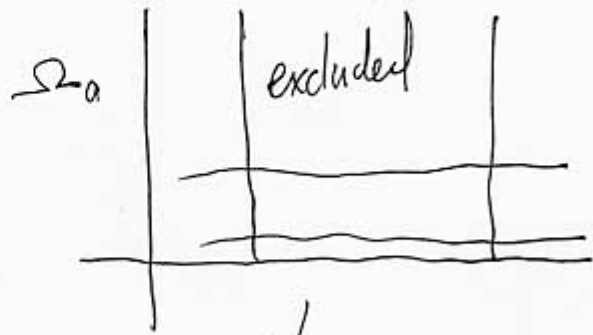
↕ analogous interaction

$$g/f \frac{c}{64\pi} \frac{F_{\mu\nu} F_{\rho\sigma} \epsilon^{\mu\nu\rho\sigma}}{4 \vec{E} \cdot \vec{B}}$$

one search method



$\sim \Omega_a$ but can be 0 in principle



Sketch of exclusion plot, see PDG for review details.