

HW #4

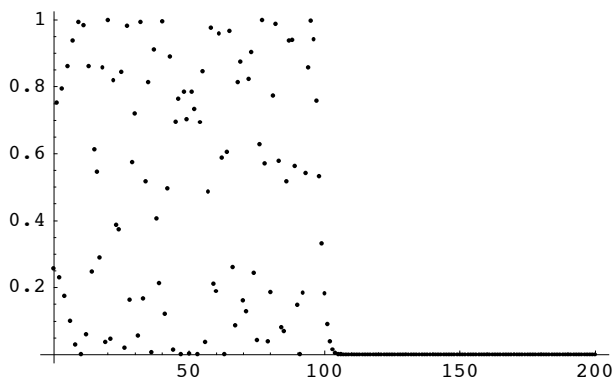
1. Hard Sphere Scattering

For the hard sphere scattering, the requirement on the radial wave function is $R_l(r) = j_l(kr) \cos \delta_l + n_l(kr) \sin \delta_l$ for $r > a$, and $R_l(a) = 0$. Therefore, $\tan \delta_l = -\frac{j_l(ka)}{n_l(ka)}$. Using the trigonometric identity $\sin^2 \theta = \frac{\tan^2 \theta}{1 + \tan^2 \theta} = \frac{1}{1 + \cos^2 \theta}$, and for the case $ka = 100$,

```
sin2delta =
```

$$\text{Table}\left[\left\{1, \text{N}\left[\frac{(\text{BesselJ}\left[\frac{(2l+1)}{2}, z\right])^2}{(\text{BesselJ}\left[\frac{(2l+1)}{2}, z\right])^2 + (\text{BesselY}\left[\frac{(2l+1)}{2}, z\right])^2} /. \{z \rightarrow 100\}\right]\right\}, \{1, 0, 200\}\right];$$

```
ListPlot[sin2delta]
```



```
- Graphics -
```

It appears that $\sin^2 \delta_l$ behaves more or less randomly between 0 and 1, and hence $\frac{1}{2}$ on average.

$$\text{Sum}\left[4\pi \frac{(2l+1)}{100^2} \text{sin2delta}[[1+l, 2]], \{1, 0, 200\}\right]$$

```
6.56891
```

We can analytically estimate the total cross section in the $ka \gg 1$ limit, by regarding $\sin^2 \delta_l = \frac{1}{2}$ on average for $0 \leq l \leq ka$. Then $\sigma = \sum_{l=0}^{ka} \frac{4\pi(2l+1)}{k^2} \frac{1}{2} = \frac{2\pi}{k^2} (ka+1)^2 = 2\pi a^2$ up to corrections of order $(ka)^{-1}$. Indeed, the total cross section above for $ka = 100$ is quite close to

$$\text{N}[2\pi a^2 /. \{a \rightarrow 1\}]$$

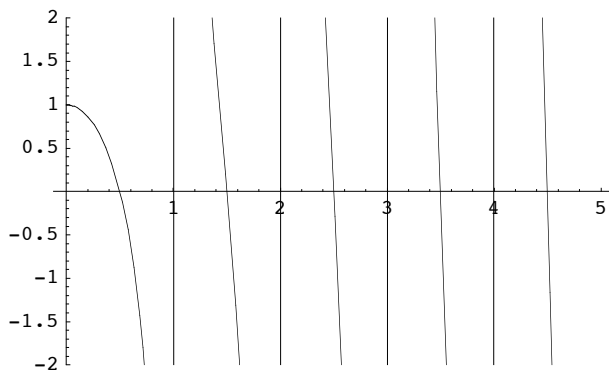
```
6.28319
```

2. Spherical Well Potential for $l = 0$. (Threshold bound states)

(a)

Bound state problem is given by $r R_0(r) = \sin \sqrt{K^2 - \kappa^2} r$ for $r < a$, and $r R_0(r) = e^{-\kappa r}$ for $r > a$. The continuity of the logarithmic derivatives requires $\sqrt{K^2 - \kappa^2} \cot \sqrt{K^2 - \kappa^2} a = -\kappa$ or $\frac{\kappa}{\sqrt{K^2 - \kappa^2}} \tan \sqrt{K^2 - \kappa^2} a + 1 = 0$. To understand where the solutions appear, we plot the function $z \cot z$

```
Plot[ $\pi x \cot[\pi x]$ , { $x$ , 0, 5}, PlotRange -> {-2, 2}]
```



- Graphics -

Note that the straight vertical lines are an artefact of the plotting by *Mathematica*.

For a bound state to form, we need $\sqrt{K^2 - \kappa^2} \cot \sqrt{K^2 - \kappa^2} a$ to be the same as $-\kappa < 0$. Therefore, the solutions arise when $(n + \frac{1}{2})\pi < \sqrt{K^2 - \kappa^2} a < (n + 1)\pi$, or $\sqrt{K^2 - ((n + 1)\pi)^2} < \kappa < \sqrt{K^2 - ((n + \frac{1}{2})\pi)^2}$.

For the threshold bound states, $\kappa = 0$ and hence $K \cot K a = 0$, and therefore $K a = (n + \frac{1}{2})\pi$. New solutions appear there, and the binding energies increase as K is increased.

The rest is not required, but this is how one can plot the binding energies as a function of the potential depth.

```

soll = Table[{K,  $\kappa$  /. FindRoot[ $\sqrt{\mathbf{K}^2 - \kappa^2} \cot[\sqrt{\mathbf{K}^2 - \kappa^2} \mathbf{a}] + \kappa$  /. {a → 1},
  { $\kappa$ ,  $\sqrt{\mathbf{K}^2 - \left(\frac{1}{2} \pi\right)^2}$ ,  $\sqrt{\mathbf{Max}[\mathbf{K}^2 - \pi^2, 0]}$ ,  $\sqrt{\mathbf{K}^2 - \left(\frac{1}{2} \pi\right)^2}$ ]}][[1]], {K,  $\frac{1}{2} \pi$ , 5  $\pi$ ,  $\frac{1}{10} \pi$ ]}

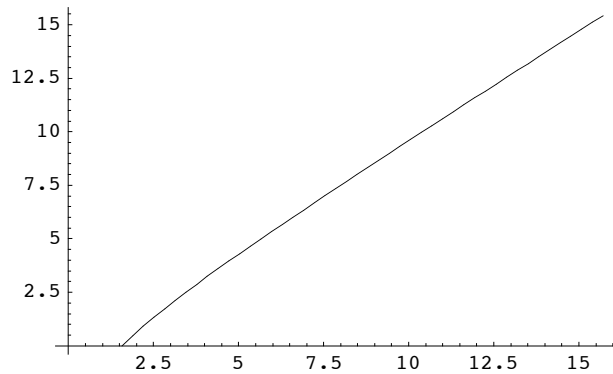
```

$\left\{ \left\{ \frac{\pi}{2}, 0. \right\}, \left\{ \frac{3\pi}{5}, 0.473252 \right\}, \left\{ \frac{7\pi}{10}, 0.914789 \right\}, \left\{ \frac{4\pi}{5}, 1.33373 \right\}, \left\{ \frac{9\pi}{10}, 1.73588 \right\}, \left\{ \pi, 2.12514 \right\}, \right.$
 $\left\{ \frac{11\pi}{10}, 2.50428 \right\}, \left\{ \frac{6\pi}{5}, 2.87526 \right\}, \left\{ \frac{13\pi}{10}, 3.2396 \right\}, \left\{ \frac{7\pi}{5}, 3.59842 \right\}, \left\{ \frac{3\pi}{2}, 3.95262 \right\}, \right.$
 $\left\{ \frac{8\pi}{5}, 4.30288 \right\}, \left\{ \frac{17\pi}{10}, 4.64979 \right\}, \left\{ \frac{9\pi}{5}, 4.99379 \right\}, \left\{ \frac{19\pi}{10}, 5.33527 \right\}, \left\{ 2\pi, 5.67453 \right\}, \right.$
 $\left\{ \frac{21\pi}{10}, 6.01184 \right\}, \left\{ \frac{11\pi}{5}, 6.34743 \right\}, \left\{ \frac{23\pi}{10}, 6.68148 \right\}, \left\{ \frac{12\pi}{5}, 7.01415 \right\}, \left\{ \frac{5\pi}{2}, 7.34559 \right\}, \right.$
 $\left\{ \frac{13\pi}{5}, 7.67591 \right\}, \left\{ \frac{27\pi}{10}, 8.00522 \right\}, \left\{ \frac{14\pi}{5}, 8.33362 \right\}, \left\{ \frac{29\pi}{10}, 8.66118 \right\}, \left\{ 3\pi, 8.98798 \right\}, \right.$
 $\left\{ \frac{31\pi}{10}, 9.31408 \right\}, \left\{ \frac{16\pi}{5}, 9.63954 \right\}, \left\{ \frac{33\pi}{10}, 9.96441 \right\}, \left\{ \frac{17\pi}{5}, 10.2887 \right\}, \left\{ \frac{7\pi}{2}, 10.6126 \right\}, \right.$
 $\left\{ \frac{18\pi}{5}, 10.9359 \right\}, \left\{ \frac{37\pi}{10}, 11.2588 \right\}, \left\{ \frac{19\pi}{5}, 11.5813 \right\}, \left\{ \frac{39\pi}{10}, 11.9035 \right\}, \left\{ 4\pi, 12.2253 \right\}, \right.$
 $\left\{ \frac{41\pi}{10}, 12.5467 \right\}, \left\{ \frac{21\pi}{5}, 12.8679 \right\}, \left\{ \frac{43\pi}{10}, 13.1887 \right\}, \left\{ \frac{22\pi}{5}, 13.5093 \right\}, \left\{ \frac{9\pi}{2}, 13.8296 \right\}, \right.$
 $\left\{ \frac{23\pi}{5}, 14.1497 \right\}, \left\{ \frac{47\pi}{10}, 14.4696 \right\}, \left\{ \frac{24\pi}{5}, 14.7893 \right\}, \left\{ \frac{49\pi}{10}, 15.1087 \right\}, \left\{ 5\pi, 15.428 \right\} \right\}$

```

plot1 = ListPlot[soll, PlotJoined → True]

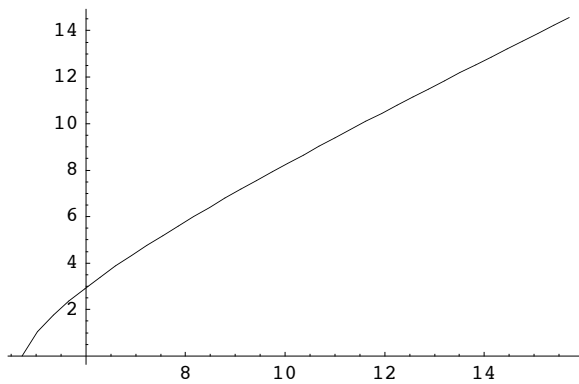
```



- Graphics -

```
sol2 = Table[{K,  $\kappa$  /. FindRoot[ $\sqrt{K^2 - \kappa^2} \text{Cot}[\sqrt{K^2 - \kappa^2} a] + \kappa /. \{a \rightarrow 1\}$ ,
  { $\kappa$ ,  $\sqrt{K^2 - \left(\frac{3}{2} \pi\right)^2}$ ,  $\sqrt{\text{Max}[K^2 - (2 \pi)^2, 0]}$ ,  $\sqrt{K^2 - \left(\frac{3}{2} \pi\right)^2}$ ]}][[1]]}, {K,  $\frac{3}{2} \pi$ ,  $5 \pi$ ,  $\frac{1}{10} \pi$ ]}
{{ $\frac{3 \pi}{2}$ , 0.}, { $\frac{8 \pi}{5}$ , 1.03272}, { $\frac{17 \pi}{10}$ , 1.74932}, { $\frac{9 \pi}{5}$ , 2.35381}, { $\frac{19 \pi}{10}$ , 2.89665},
{2  $\pi$ , 3.39955}, { $\frac{21 \pi}{10}$ , 3.87416}, { $\frac{11 \pi}{5}$ , 4.32756}, { $\frac{23 \pi}{10}$ , 4.7644}, { $\frac{12 \pi}{5}$ , 5.18794},
{ $\frac{5 \pi}{2}$ , 5.60053}, { $\frac{13 \pi}{5}$ , 6.00397}, { $\frac{27 \pi}{10}$ , 6.39963}, { $\frac{14 \pi}{5}$ , 6.78861},
{ $\frac{29 \pi}{10}$ , 7.17178}, {3  $\pi$ , 7.54986}, { $\frac{31 \pi}{10}$ , 7.92344}, { $\frac{16 \pi}{5}$ , 8.29304},
{ $\frac{33 \pi}{10}$ , 8.65906}, { $\frac{17 \pi}{5}$ , 9.02187}, { $\frac{7 \pi}{2}$ , 9.38178}, { $\frac{18 \pi}{5}$ , 9.73905},
{ $\frac{37 \pi}{10}$ , 10.0939}, { $\frac{19 \pi}{5}$ , 10.4466}, { $\frac{39 \pi}{10}$ , 10.7973}, {4  $\pi$ , 11.1462}, { $\frac{41 \pi}{10}$ , 11.4933},
{ $\frac{21 \pi}{5}$ , 11.8389}, { $\frac{43 \pi}{10}$ , 12.183}, { $\frac{22 \pi}{5}$ , 12.5257}, { $\frac{9 \pi}{2}$ , 12.8672},
{ $\frac{23 \pi}{5}$ , 13.2076}, { $\frac{47 \pi}{10}$ , 13.5468}, { $\frac{24 \pi}{5}$ , 13.885}, { $\frac{49 \pi}{10}$ , 14.2222}, {5  $\pi$ , 14.5585}}
```

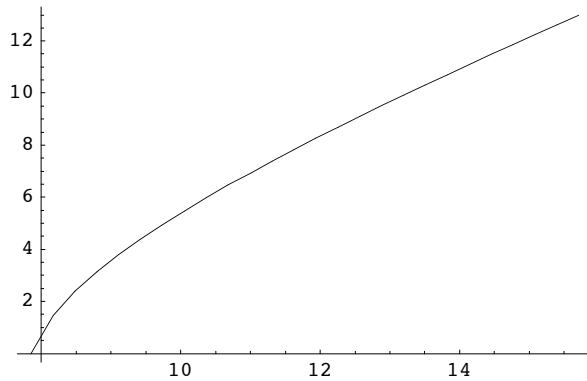
```
plot2 = ListPlot[sol2, PlotJoined  $\rightarrow$  True]
```



- Graphics -

```
sol3 = Table[{K,  $\kappa$  /. FindRoot[ $\sqrt{K^2 - \kappa^2} \text{Cot}[\sqrt{K^2 - \kappa^2} a] + \kappa /. \{a \rightarrow 1\}$ ,
  { $\kappa$ ,  $\sqrt{K^2 - \left(\frac{5}{2} \pi\right)^2}$ ,  $\sqrt{\text{Max}[K^2 - (3 \pi)^2, 0]}$ ,  $\sqrt{K^2 - \left(\frac{5}{2} \pi\right)^2}$ ]}][[1]]}, {K,  $\frac{5}{2} \pi$ ,  $5 \pi$ ,  $\frac{1}{10} \pi$ ]}
{{ $\frac{5 \pi}{2}$ , 0.}, { $\frac{13 \pi}{5}$ , 1.46948}, { $\frac{27 \pi}{10}$ , 2.38764}, { $\frac{14 \pi}{5}$ , 3.13592},
{ $\frac{29 \pi}{10}$ , 3.79322}, {3  $\pi$ , 4.39218}, { $\frac{31 \pi}{10}$ , 4.95}, { $\frac{16 \pi}{5}$ , 5.477}, { $\frac{33 \pi}{10}$ , 5.97992},
{ $\frac{17 \pi}{5}$ , 6.46347}, { $\frac{7 \pi}{2}$ , 6.93106}, { $\frac{18 \pi}{5}$ , 7.38525}, { $\frac{37 \pi}{10}$ , 7.82803},
{ $\frac{19 \pi}{5}$ , 8.26096}, { $\frac{39 \pi}{10}$ , 8.6853}, {4  $\pi$ , 9.1021}, { $\frac{41 \pi}{10}$ , 9.5122},
{ $\frac{21 \pi}{5}$ , 9.91633}, { $\frac{43 \pi}{10}$ , 10.3151}, { $\frac{22 \pi}{5}$ , 10.709}, { $\frac{9 \pi}{2}$ , 11.0986},
{ $\frac{23 \pi}{5}$ , 11.4841}, { $\frac{47 \pi}{10}$ , 11.866}, { $\frac{24 \pi}{5}$ , 12.2445}, { $\frac{49 \pi}{10}$ , 12.62}, {5  $\pi$ , 12.9926}}
```

```
plot3 = ListPlot[sol3, PlotJoined → True]
```



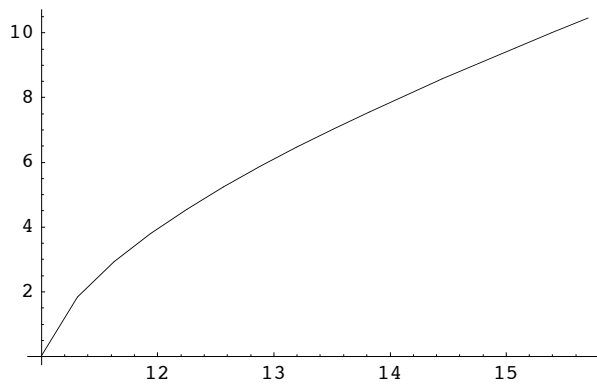
- Graphics -

```
sol4 = Table[{K, κ /. FindRoot[√(K² - κ²) Cot[√(K² - κ²) a] + κ /. {a → 1},
  {κ, √(K² - (7/2 π)²), √(Max[K² - (4 π)², 0]), √(K² - (7/2 π)²)}][[1]]}, {K, 7/2 π, 5 π, 1/10 π}]
```

FindRoot::reged : The point {0.} is at the edge of the search region {0., 0.} in coordinate 1 and the computed search direction points outside the region. **More...**

```
{{{7/2 π, 0.}, {18/5 π, 1.84024}, {37/10 π, 2.92474}, {19/5 π, 3.79432},
  {39/10 π, 4.54997}, {4 π, 5.23274}, {41/10 π, 5.86408}, {21/5 π, 6.45683},
  {43/10 π, 7.01939}, {22/5 π, 7.55759}, {9/2 π, 8.07566}, {23/5 π, 8.57679},
  {47/10 π, 9.06345}, {24/5 π, 9.53759}, {49/10 π, 10.0008}, {5 π, 10.4543}}
```

```
plot4 = ListPlot[sol4, PlotJoined → True]
```



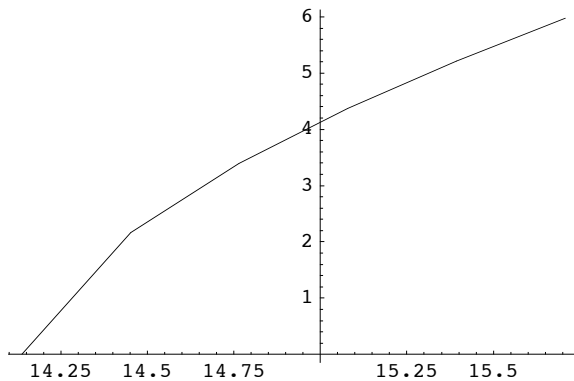
- Graphics -

```
sol5 = Table[{K, κ /. FindRoot[ $\sqrt{K^2 - \kappa^2} \cot[\sqrt{K^2 - \kappa^2} a] + \kappa /. {a \rightarrow 1},$ 
  {κ,  $\sqrt{K^2 - \left(\frac{9}{2} \pi\right)^2}$ ,  $\sqrt{\text{Max}[K^2 - (5 \pi)^2, 0]}$ ,  $\sqrt{K^2 - \left(\frac{9}{2} \pi\right)^2}$ ]}][[1]]}, {K,  $\frac{9}{2} \pi$ , 5 π,  $\frac{1}{10} \pi$ }]
```

```
FindRoot::reged : The point {0.} is at the edge of the search region {0., 0.}
  in coordinate 1 and the computed search direction points outside the region. More...
```

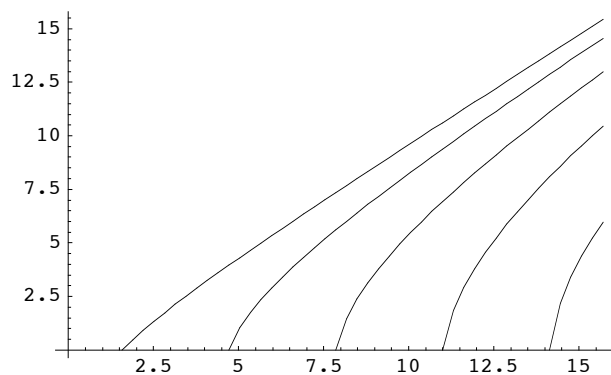
```
{ $\left\{\frac{9 \pi}{2}, 0.\right\}$ ,  $\left\{\frac{23 \pi}{5}, 2.1681\right\}$ ,  $\left\{\frac{47 \pi}{10}, 3.39733\right\}$ ,
  $\left\{\frac{24 \pi}{5}, 4.37371\right\}$ ,  $\left\{\frac{49 \pi}{10}, 5.21678\right\}$ ,  $\{5 \pi, 5.97463\}$ }
```

```
plot5 = ListPlot[sol5, PlotJoined → True]
```



- Graphics -

```
Show[plot1, plot2, plot3, plot4, plot5]
```



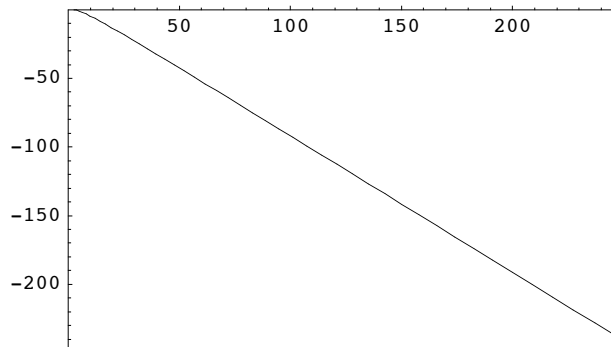
- Graphics -

This plot shows the relationship between K on the horizontal axis and κ on the vertical axis. New bound states indeed appear at $K = (n + \frac{1}{2}) \frac{\pi}{a}$. To plot the bound state energies $E = -\hbar^2 \kappa^2 / 2m$ as a function of $V_0 = 2mK^2 / \hbar^2$, we need to just square them in the unit of $2m / \hbar^2 = 1$.

```
Dimensions[sol1]
```

```
{46, 2}
```

```
E1 = ListPlot[Table[{sol1[[i, 1]]2, -sol1[[i, 2]]2}, {i, 1, 46}],  
PlotJoined → True, PlotRange → {{0, (5 π)2}, {-(5 π)2, 0}}]
```

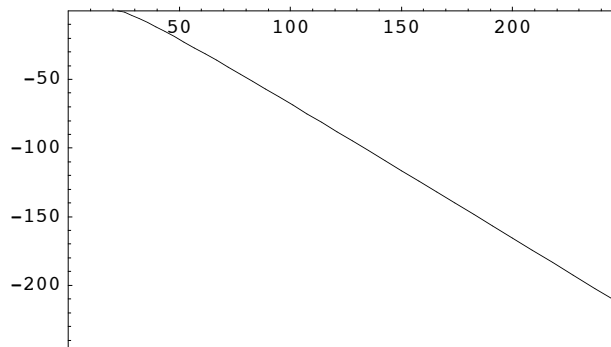


- Graphics -

```
Dimensions[sol2]
```

```
{36, 2}
```

```
E2 = ListPlot[Table[{sol2[[i, 1]]2, -sol2[[i, 2]]2}, {i, 1, 36}],  
PlotJoined → True, PlotRange → {{0, (5 π)2}, {-(5 π)2, 0}}]
```

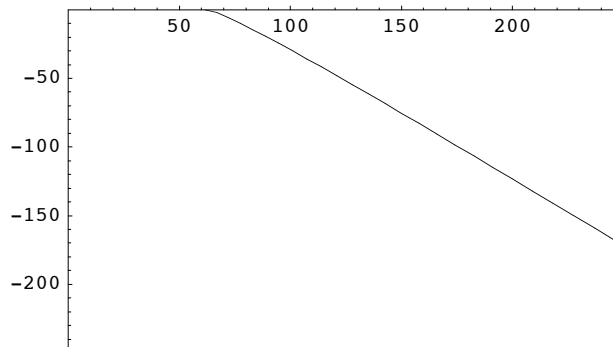


- Graphics -

```
Dimensions[sol3]
```

```
{26, 2}
```

```
E3 = ListPlot[Table[{sol3[[i, 1]]2, -sol3[[i, 2]]2}, {i, 1, 26}],  
PlotJoined → True, PlotRange → {{0, (5 π)2}, {-(5 π)2, 0}}]
```

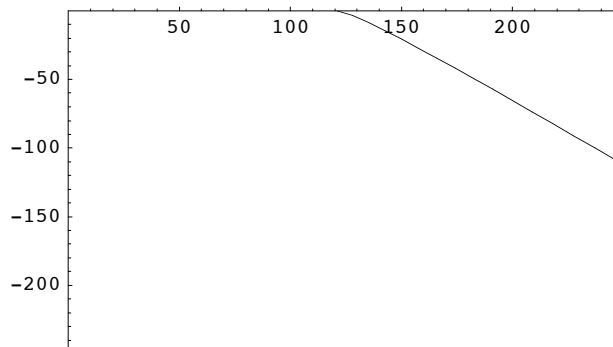


- Graphics -

```
Dimensions[sol4]
```

```
{16, 2}
```

```
E4 = ListPlot[Table[{sol4[[i, 1]]2, -sol4[[i, 2]]2}, {i, 1, 16}],  
PlotJoined → True, PlotRange → {{0, (5 π)2}, {-(5 π)2, 0}}]
```



- Graphics -

```
Dimensions[sol5]
```

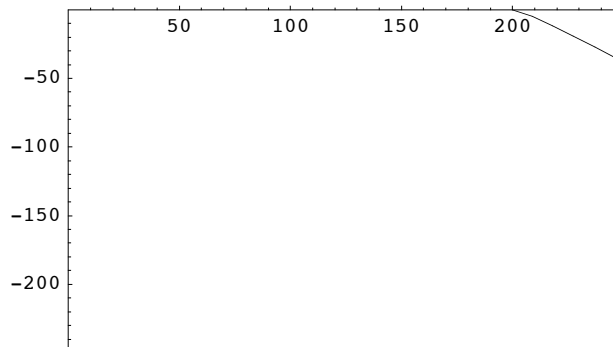
```
{6, 2}
```



```

E5 = ListPlot[Table[{sol5[[i, 1]]^2, -sol5[[i, 2]]^2}, {i, 1, 6}],
  PlotJoined -> True, PlotRange -> {{0, (5 π)^2}, {-(5 π)^2, 0}}]

```

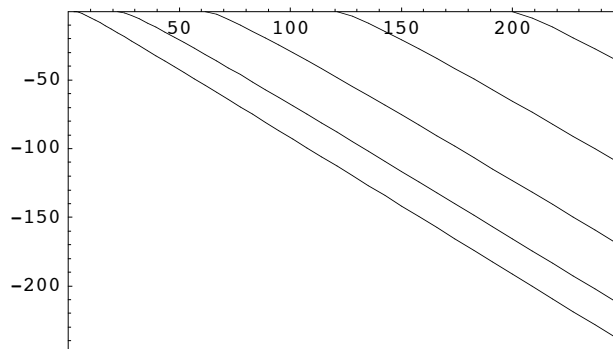


- Graphics -

```

Show[E1, E2, E3, E4, E5]

```



- Graphics -

This plot shows how new bound states appear at each $V_0 = \frac{1}{2ma^2} \hbar^2 \left(n + \frac{1}{2}\right)^2 \pi^2$ as the depth of the potential is increased.

(b)

Phase shift is given by matching $r R_0(r) = \sin \sqrt{K^2 + k^2} r$ for $r < a$ and $r R_0(r) = \sin k r \cos \delta_0 + \cos k r \sin \delta_0 = \sin(k r + \delta_0)$. The logarithmic derivatives are $\sqrt{K^2 + k^2} \cot \sqrt{K^2 + k^2} a = k \cot(k a - \delta_0)$, or $\frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a = \tan(k a + \delta_0)$. therefore, $\delta_0 = -k a + \arctan \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a$.

We can rewrite it as

$$\begin{aligned} \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a &= \tan(k a + \delta_0) \\ &= -i \frac{e^{i(k a + \delta_0)} - e^{-i(k a + \delta_0)}}{e^{i(k a + \delta_0)} + e^{-i(k a + \delta_0)}} = -i \frac{e^{i k a} e^{2 i \delta_0} - e^{-i k a}}{e^{i k a} e^{2 i \delta_0} + e^{-i k a}} \end{aligned}$$

we find

$$i \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a (e^{i k a} e^{2 i \delta_0} + e^{-i k a}) = e^{i k a} e^{2 i \delta_0} - e^{-i k a}$$

and

$$e^{2 i \delta_0} e^{i k a} \left(i \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a - 1 \right) = \left(-i \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a - 1 \right)$$

and hence

$$e^{2 i \delta_0} = e^{-i k a} \frac{1 + i \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a}{1 - i \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a}$$

(c)

The S -matrix element $e^{2 i \delta_0}$ has poles when

$$1 - i \frac{k}{\sqrt{K^2 + k^2}} \tan \sqrt{K^2 + k^2} a = 0.$$

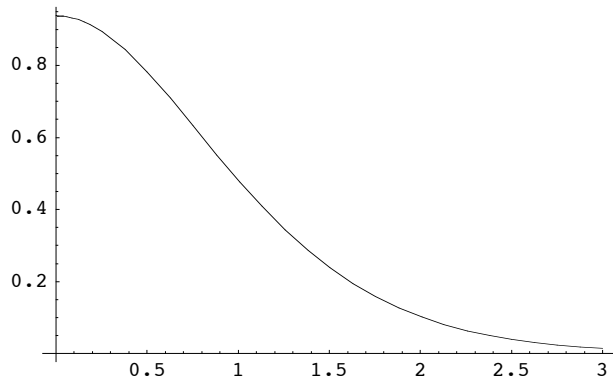
On the upper half plane, we can write $k = i \kappa$, where $\text{Re}(\kappa) > 0$. The poles appear where

$1 + \frac{\kappa}{\sqrt{K^2 - \kappa^2}} \tan \sqrt{K^2 - \kappa^2} a = 0$. This is precisely the condition for the bound states studied in part (a).

$$\begin{aligned} \delta_0 &= \text{ArcTan} \left[\frac{\mathbf{k}}{\sqrt{\mathbf{k}^2 + \mathbf{K}^2}} \text{Tan} \left[\sqrt{\mathbf{k}^2 + \mathbf{K}^2} \mathbf{a} \right] \right] - \mathbf{k} \mathbf{a} \\ &= -\mathbf{a} \mathbf{k} + \text{ArcTan} \left[\frac{\mathbf{k} \text{Tan} \left[\mathbf{a} \sqrt{\mathbf{k}^2 + \mathbf{K}^2} \right]}{\sqrt{\mathbf{k}^2 + \mathbf{K}^2}} \right] \end{aligned}$$

Far away from a threshold bound state, for instance for $K = \frac{\pi}{4}$,

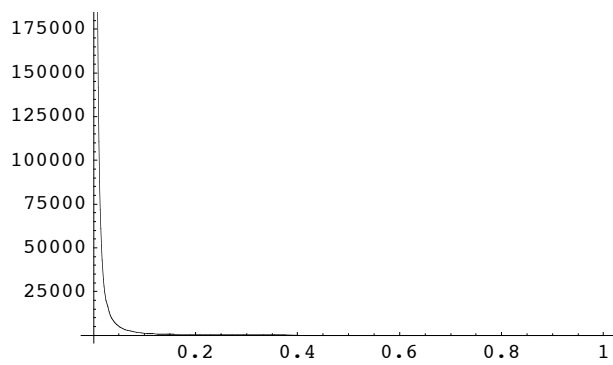
Plot $\left[\frac{4 \pi}{k^2} \text{Sin}[\delta_0]^2 /. \{a \rightarrow 1, \kappa \rightarrow \frac{\pi}{4}\}, \{k, 0, 3\} \right]$



- Graphics -

which is smaller than the geometric cross section. On the other hand, exactly on the threshold resonance,

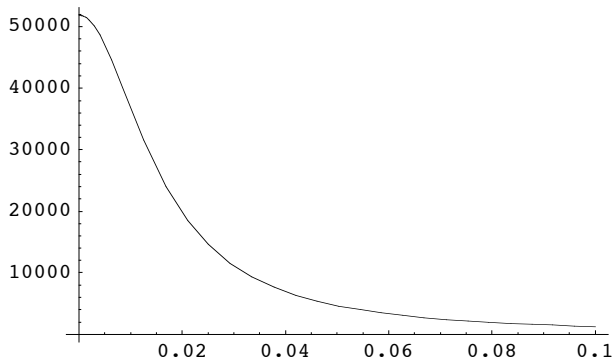
Plot $\left[\frac{4 \pi}{k^2} \text{Sin}[\delta_0]^2 /. \{a \rightarrow 1, \kappa \rightarrow \frac{\pi}{2}\}, \{k, 0, 1\} \right]$



- Graphics -

and the cross section diverges at $k = 0$. For K just above the threshold bound state, namely when the bound state exists just below $E = 0$,

`Plot[$\frac{4 \pi}{k^2} \text{Sin}[\delta_0]^2 /. \{a \rightarrow 1, K \rightarrow \frac{\pi}{2} + 0.01\}, \{k, 0, 0.1\}]$`

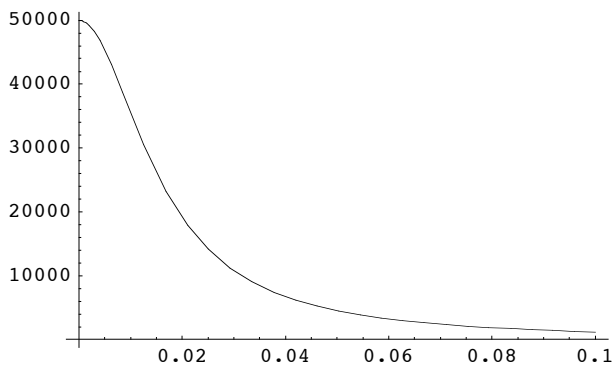


- Graphics -

The cross section is finite, but is much bigger than the geometric cross section. Similarly, for K just below the threshold bound state, namely when the bound state has just disappeared above $E = 0$,

- Graphics -

`Plot[$\frac{4 \pi}{k^2} \text{Sin}[\delta_0]^2 /. \{a \rightarrow 1, K \rightarrow \frac{\pi}{2} - 0.01\}, \{k, 0, 0.1\}]$`



- Graphics -

It still exhibits a cross section much larger than the geometric cross section.

3. Spherical Well Potential for $l = 1$. (Resonances)

(a)

Here is $j_1(z)$:

$$\text{PowerExpand}\left[\sqrt{\frac{\pi}{2z}} \text{BesselJ}\left[\frac{3}{2}, z\right]\right]$$

$$\frac{-\text{Cos}[z] + \frac{\text{Sin}[z]}{z}}{z}$$

Here is $n_1(z)$:

$$\text{PowerExpand}\left[\sqrt{\frac{\pi}{2z}} \text{BesselY}\left[\frac{3}{2}, z\right]\right]$$

$$\frac{-\frac{\text{Cos}[z]}{z} - \text{Sin}[z]}{z}$$

We use $r R_1(r)$

$$\text{Rin} = -\text{Cos}\left[\sqrt{k^2 + K^2} r\right] + \frac{\text{Sin}\left[\sqrt{k^2 + K^2} r\right]}{\sqrt{k^2 + K^2} r}$$

$$-\text{Cos}\left[\sqrt{k^2 + K^2} r\right] + \frac{\text{Sin}\left[\sqrt{k^2 + K^2} r\right]}{\sqrt{k^2 + K^2} r}$$

$$\text{Rout} = \left(-\text{Cos}[kr] + \frac{\text{Sin}[kr]}{kr}\right) \text{Cos}[\delta_1] + \left(\frac{-\text{Cos}[kr]}{kr} - \text{Sin}[kr]\right) \text{Sin}[\delta_1]$$

General::spell1 : Possible spelling error: new symbol name "Rout" is similar to existing symbol "Root". MORE...

$$\text{Cos}[\delta_1] \left(-\text{Cos}[kr] + \frac{\text{Sin}[kr]}{kr}\right) + \left(-\frac{\text{Cos}[kr]}{kr} - \text{Sin}[kr]\right) \text{Sin}[\delta_1]$$

Simplify[Rout]

$$-\text{Cos}[kr - \delta_1] + \frac{\text{Sin}[kr - \delta_1]}{kr}$$

Simplify $\left[\frac{\text{D}[\text{Rin}, r]}{\text{Rin}}\right]$

$$-\left(\sqrt{k^2 + K^2} r \text{Cos}\left[\sqrt{k^2 + K^2} r\right] + (-1 + k^2 r^2 + K^2 r^2) \text{Sin}\left[\sqrt{k^2 + K^2} r\right]\right) /$$

$$\left(r \left(\sqrt{k^2 + K^2} r \text{Cos}\left[\sqrt{k^2 + K^2} r\right] - \text{Sin}\left[\sqrt{k^2 + K^2} r\right]\right)\right)$$

D[Rout, r]

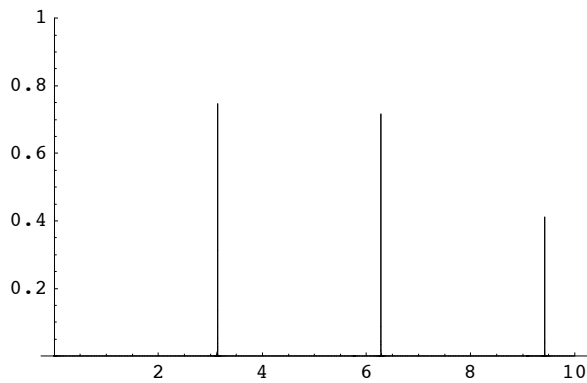
Rout

$$\left(\text{Cos}[\delta_1] \left(\frac{\text{Cos}[kr]}{r} + k \text{Sin}[kr] - \frac{\text{Sin}[kr]}{kr^2}\right) + \left(-k \text{Cos}[kr] + \frac{\text{Cos}[kr]}{kr^2} + \frac{\text{Sin}[kr]}{r}\right) \text{Sin}[\delta_1]\right) /$$

$$\left(\text{Cos}[\delta_1] \left(-\text{Cos}[kr] + \frac{\text{Sin}[kr]}{kr}\right) + \left(-\frac{\text{Cos}[kr]}{kr} - \text{Sin}[kr]\right) \text{Sin}[\delta_1]\right)$$

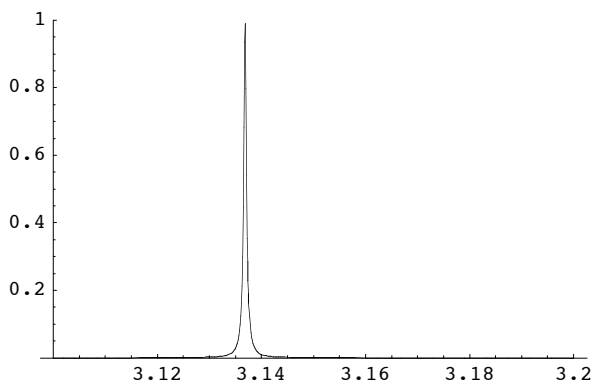
```
sol = Solve[-(sqrt[k^2 + K^2] r Cos[sqrt[k^2 + K^2] r] + (-1 + k^2 r^2 + K^2 r^2) Sin[sqrt[k^2 + K^2] r]) /
  (r (sqrt[k^2 + K^2] r Cos[sqrt[k^2 + K^2] r] - Sin[sqrt[k^2 + K^2] r])) ==
  (cot (Cos[k r] / r + k Sin[k r] - Sin[k r] / (k r^2)) + (-k Cos[k r] + Cos[k r] / (k r^2) + Sin[k r] / r)) /
  (cot (-Cos[k r] + Sin[k r] / (k r)) + (-Cos[k r] / (k r) - Sin[k r])) /. {r -> a}, cot]
{{cot -> (-a k^2 sqrt[k^2 + K^2] Cos[a k] Cos[a sqrt[k^2 + K^2]] - K^2 Cos[a k] Sin[a sqrt[k^2 + K^2]] -
  a k^3 Sin[a k] Sin[a sqrt[k^2 + K^2]] - a k K^2 Sin[a k] Sin[a sqrt[k^2 + K^2]]) /
  (-a k^2 sqrt[k^2 + K^2] Cos[a sqrt[k^2 + K^2]] Sin[a k] + a k^3 Cos[a k] Sin[a sqrt[k^2 + K^2]] +
  a k K^2 Cos[a k] Sin[a sqrt[k^2 + K^2]] - K^2 Sin[a k] Sin[a sqrt[k^2 + K^2]])}}
```

```
Plot[1 / (1 + cot^2) /. sol[[1]] /. {a -> 1} /. {k -> 0.1},
  {K, 0, 10}, PlotRange -> {0, 1}, PlotPoints -> 500]
```



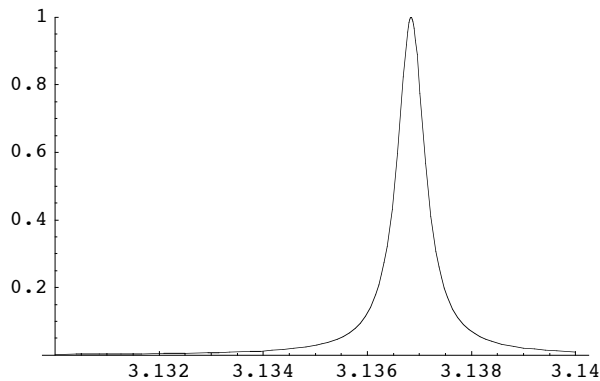
- Graphics -

```
Plot[1 / (1 + cot^2) /. sol[[1]] /. {a -> 1} /. {k -> 0.1}, {K, 3.1, 3.2}, PlotRange -> {0, 1}]
```



- Graphics -

```
Plot[ $\frac{1}{1 + \cot^2}$  /. sol[[1]] /. {a -> 1} /. {k -> 0.1}, {K, 3.13, 3.14}, PlotRange -> {0, 1}]
```



- Graphics -

```
firstpeak = FindRoot[cot /. sol[[1]] /. {a -> 1} /. {k -> 0.1}, {K, 3.137}]
```

```
{K -> 3.13684}
```

The corresponding value of V_0 at the peak is

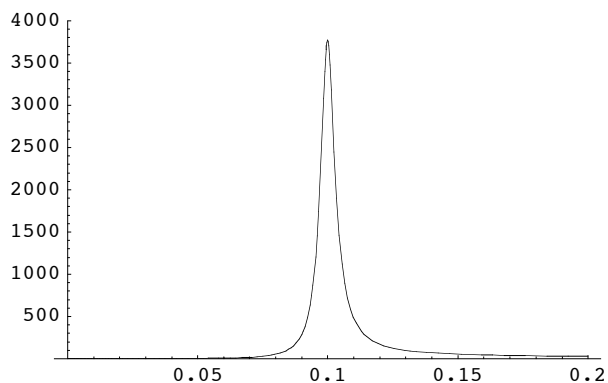
$$\frac{\hbar^2 K^2}{2 m a^2} /. \text{firstpeak}$$

$$\frac{4.91989 \hbar^2}{a^2 m}$$

(b)

The cross section is dominated by the $l = 1$ partial wave around the peak.

```
Plot[ $\frac{4 \pi}{k^2} 3 \frac{1}{1 + \cot^2}$  /. sol[[1]] /. {a -> 1} /. firstpeak, {k, 0, 0.2}, PlotRange -> {0, 4000}]
```



- Graphics -

It shows a prominent peak at $k a = 0.1$, and the cross section is much larger than the geometric cross section $4 \pi a^2$.

(C)

We need a precise value for the phase shift at the peak

```
sol[[1]] /. {a -> 1, k -> 0.1} /. firstpeak
```

```
{cot -> -2.81908 × 10-13}
```

```
 $\delta_1 = \text{ArcTan}\left[\frac{1}{\text{cot}}\right] /. \%$ 
```

```
-1.5708
```

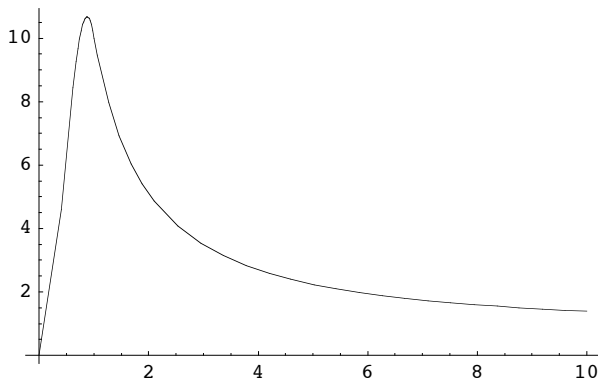
```
Rout /. { $\delta_1 \rightarrow -1.5707963267946148^\wedge$ } /. firstpeak /. {a -> 1, k -> 0.1} /. {r -> 1}
```

```
10.0499
```

```
Rin /. { $\delta_1 \rightarrow -1.5707963267946148^\wedge$ } /. firstpeak /. {a -> 1, k -> 0.1} /. {r -> 1}
```

```
1.001
```

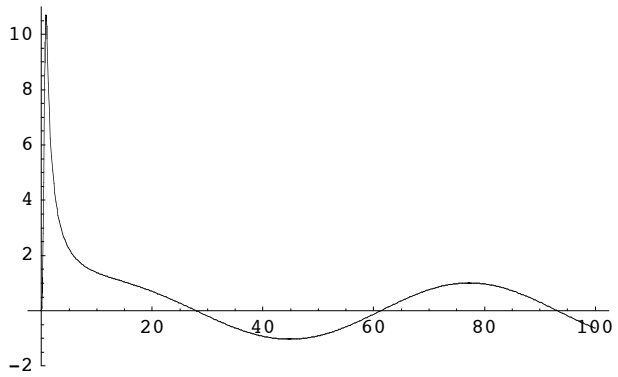
```
Plot[If[r > 1, Rout /. sol[[1]] /. { $\delta_1 \rightarrow -1.5707963267946148^\wedge$ },  $\frac{10.049875069427088^\wedge}{1.0010011782275658^\wedge}$  Rin] /.  
firstpeak /. {a -> 1, k -> 0.1}, {r, 0, 10}]
```



- Graphics -

It appears more-or-less a bound-state solution at small r , because the classical turning point is $\frac{l(l+1)}{r^2} = k^2$ and hence $r = \frac{\sqrt{2}}{k} \approx 14$ in our case. However, it does oscillate far out in radius as expected for a non-bound state beyond the classical turning point.


```
Plot[If[r > 1, Rout /. sol[[1]] /. { $\delta_1 \rightarrow -1.5707963267946148$ },  $\frac{10.049875069427088}{1.0010011782275658}$  Rin] /.  
  firstpeak /. {a → 1, k → 0.1}, {r, 0, 100}, PlotRange → {-2, 11}, PlotPoints → 500]
```



- Graphics -