

# HW #1

## 1. Free-particle Schrödinger Equation

(1) Plane wave  $\psi = e^{ikz}$  does not depend on  $x$  or  $y$ . Therefore, the Schrödinger equation becomes  $(\partial_z^2 + k^2)\psi = 0$ . Obviously this is a solution to the equation.

$$\mathbf{D}[\mathbf{E}^{ikz}, \{\mathbf{z}, 2\}] + \mathbf{k}^2 \mathbf{E}^{ikz}$$

0

(2) In polar coordinates, the Laplacian can be rewritten as  $\nabla^2 = \partial_r^2 + \frac{2}{r} \partial_r + \frac{1}{r^2} \partial_\theta^2 + \frac{\cos \theta}{r^2 \sin \theta} \partial_\theta + \frac{1}{r^2 \sin^2 \theta} \partial_\phi^2$ . The spherical wave  $\psi = \frac{e^{ikr}}{r}$  does not depend on  $\theta$  or  $\phi$ . Therefore, the Schrödinger equation becomes  $(\partial_r^2 + \frac{2}{r} \partial_r + k^2)\psi = 0$ .

$$\mathbf{D}\left[\frac{\mathbf{E}^{ikr}}{\mathbf{r}}, \{\mathbf{r}, 2\}\right] + \frac{2}{\mathbf{r}} \mathbf{D}\left[\frac{\mathbf{E}^{ikr}}{\mathbf{r}}, \mathbf{r}\right] + \mathbf{k}^2 \frac{\mathbf{E}^{ikr}}{\mathbf{r}}$$

$$\frac{2 e^{ikr}}{r^3} - \frac{2 i e^{ikr} k}{r^2} + \frac{2 \left(-\frac{e^{ikr}}{r^2} + \frac{i e^{ikr} k}{r}\right)}{r}$$

$$\mathbf{Simplify}[\%]$$

0

## 2. Double Pin-hole Experiment

(1) As directed, we assume that the denominators are approximately the same between two waves. This is justified because the corrections are only of the order of  $d/L$ , and we are interested in the case where  $d \ll L$ . We require that the numerators have the same phase, namely  $kr_+ - kr_- = 2\pi n$ . We expand the l.h.s. with respect to  $d$ ,

$$\mathbf{Series}[\mathbf{Sqrt}[\mathbf{x}^2 + \left(\mathbf{y} + \frac{\mathbf{d}}{2}\right)^2 + \mathbf{L}^2], \{\mathbf{d}, \mathbf{0}, \mathbf{1}\}]$$

$$\sqrt{\mathbf{L}^2 + \mathbf{x}^2 + \mathbf{y}^2} + \frac{\mathbf{y} \mathbf{d}}{2 \sqrt{\mathbf{L}^2 + \mathbf{x}^2 + \mathbf{y}^2}} + \mathcal{O}[\mathbf{d}]^2$$

$$\mathbf{Series}[\mathbf{Sqrt}[\mathbf{x}^2 + \left(\mathbf{y} - \frac{\mathbf{d}}{2}\right)^2 + \mathbf{L}^2], \{\mathbf{d}, \mathbf{0}, \mathbf{1}\}]$$

$$\sqrt{\mathbf{L}^2 + \mathbf{x}^2 + \mathbf{y}^2} - \frac{\mathbf{y} \mathbf{d}}{2 \sqrt{\mathbf{L}^2 + \mathbf{x}^2 + \mathbf{y}^2}} + \mathcal{O}[\mathbf{d}]^2$$

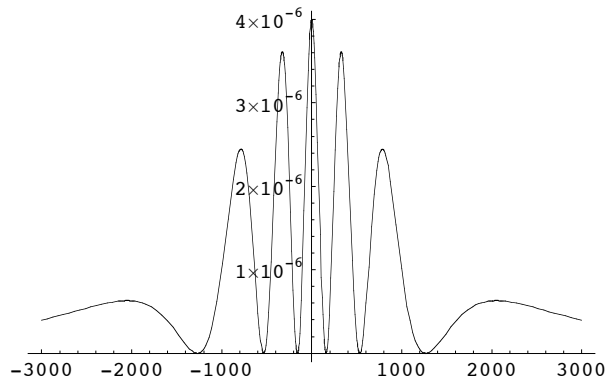
$$\mathbf{Simplify}[\mathbf{Normal}[\% - \%]]$$

$$\frac{\mathbf{d} \mathbf{y}}{\sqrt{\mathbf{L}^2 + \mathbf{x}^2 + \mathbf{y}^2}}$$

Therefore,  $k \frac{dy}{\sqrt{L^2+x^2+y^2}} = 2\pi n$  and hence  $\frac{y}{\sqrt{L^2+x^2+y^2}} = n \frac{\lambda}{d}$ ,

(2) Let us choose the unit where  $k = 1$ . Then we pick  $d = 20$ ,  $L = 1000$ . Here is the interference pattern. First along the  $y$ -axis ( $x = 0$ ):

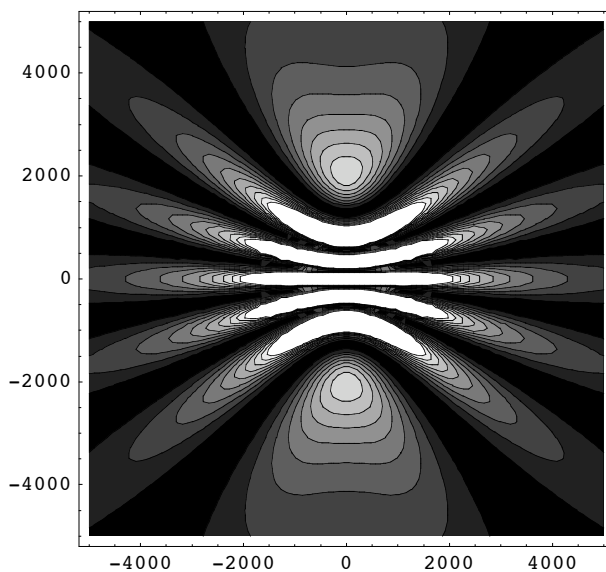
```
Plot[Abs[ $\frac{E^{i r_+}}{r_+} + \frac{E^{i r_-}}{r_-}$ ]^2 /. {r_+ ->  $\sqrt{x^2 + (y - \frac{d}{2})^2 + L^2}$ , r_- ->  $\sqrt{x^2 + (y + \frac{d}{2})^2 + L^2}$ } /.
{d -> 20, L -> 1000} /. {x -> 0}, {y, -3000, 3000}, PlotPoints -> 100]
```



- Graphics -

(3) Now on the plane:

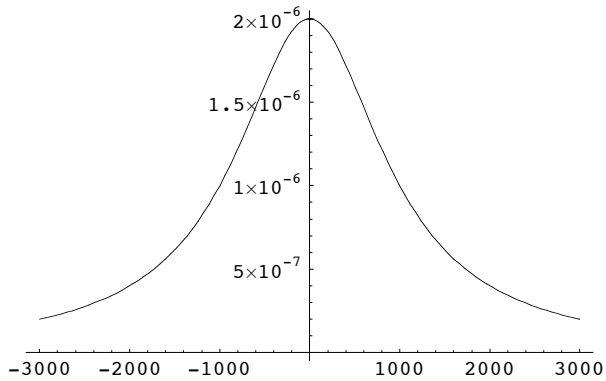
```
ContourPlot[Abs[ $\frac{E^{i r_+}}{r_+} + \frac{E^{i r_-}}{r_-}$ ]^2 /. {r_+ ->  $\sqrt{x^2 + (y - \frac{d}{2})^2 + L^2}$ , r_- ->  $\sqrt{x^2 + (y + \frac{d}{2})^2 + L^2}$ } /.
{d -> 20, L -> 1000}, {x, -5000, 5000}, {y, -5000, 5000}, PlotPoints -> 100]
```



- ContourGraphics -

(4) For the same parameters as in (2), First along the y-axis ( $x = 0$ ):

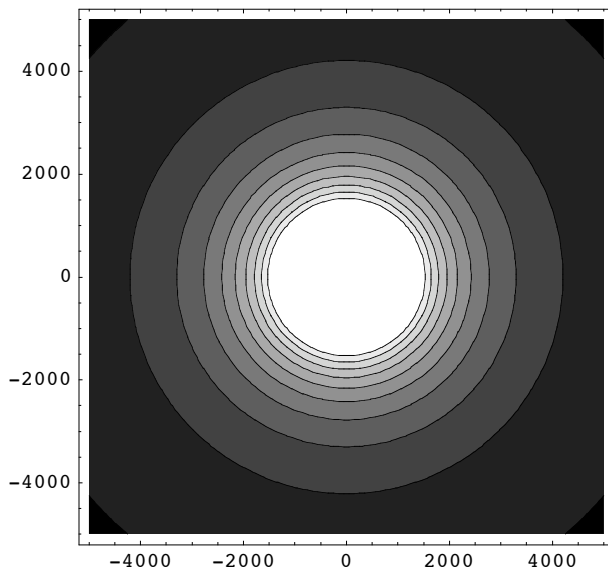
```
Plot[Abs[ $\frac{E^{i r_+}}{r_+}$ ]2 + Abs[ $\frac{E^{i r_-}}{r_-}$ ]2 /. {r+ ->  $\sqrt{x^2 + (y + \frac{d}{2})^2 + L^2}$ , r- ->  $\sqrt{x^2 + (y - \frac{d}{2})^2 + L^2}$ } /.  
{d -> 20, L -> 1000} /. {x -> 0}, {y, -3000, 3000}, PlotPoints -> 100]
```



- Graphics -

Now on the plane:

```
ContourPlot[Abs[ $\frac{E^{i r_+}}{r_+}$ ]2 + Abs[ $\frac{E^{i r_-}}{r_-}$ ]2 /. {r+ ->  $\sqrt{x^2 + (y + \frac{d}{2})^2 + L^2}$ , r- ->  $\sqrt{x^2 + (y - \frac{d}{2})^2 + L^2}$ } /.  
{d -> 20, L -> 1000}, {x, -5000, 5000}, {y, -5000, 5000}, PlotPoints -> 100]
```



- ContourGraphics -

The main difference is the absence of the interference pattern.