

Fall, 02

Physics 129A

Solutions to H/w #1

- ① Denote E = energy of the electron,
 E' = energy of the positron,
 \mathcal{E} = energy of $\Upsilon(4S)$,

p, p', p = momenta of the electron, the positron,
 and $\Upsilon(4S)$ respectively

Conservation laws imply:

$$(1) \quad \mathcal{E} = E + E',$$

$$(2) \quad p = p - p' \quad (p' \text{ is the absolute value of momentum, so we need a minus sign because the positron is going in the opposite direction of the electron})$$

We also know the masses of the three particles - denote them m for e^- and e^+ and M for $\Upsilon(4S)$

Energy and momentum of any particle are related through its mass, so we have:

$$(3), (4) \quad \frac{1}{c^2} E^2 - p^2 = m^2 c^2 = \frac{1}{c^2} E'^2 - p'^2$$

$$(5) \quad \frac{1}{c^2} \mathcal{E}^2 - p^2 = M^2 c^2$$

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We have 5 equations with five unknowns: E', E, p, p', β .

We need to solve for E' :

$$(1), (2) \rightarrow (5) \quad \frac{1}{2}(E+E')^2 - (p-p')^2 = M^2 c^2$$

$$(3), (4) \rightarrow \begin{aligned} & \frac{1}{c^2}(E^2 + 2EE' + E'^2) - p^2 + 2pp' - p'^2 = M^2 c^2 \\ & m^2 c^2 + \frac{2EE'}{c^2} + m^2 c^2 + 2pp' = M^2 c^2 \end{aligned}$$

$$(*) \quad 2c^2 pp' = (M^2 c^4 - 2m^2 c^4) - 2EE'$$

$$[A \equiv M^2 c^4 - 2m^2 c^4] \quad 4c^4 p^2 p'^2 = A^2 - 4EE'A + 4E^2 E'^2$$

$$(3), (4) \rightarrow 4(E^2 - m^2 c^4)(E'^2 - m^2 c^4) = A^2 - 4EE'A + 4E^2 E'^2$$

$$0 = 4E'^2(E^2 - m^2 c^4 - E^2) + 4EAE' - (A^2 + 4m^2 c^4(E^2 - m^2 c^4))$$

$$0 = 4m^2 c^4 E'^2 - 4EAE' + \underbrace{(M^4 c^8 + 4m^2 c^4(E^2 - M^2 c^4))}_B$$

$$E' = \frac{4EA \pm \sqrt{16E^2 A^2 - 16m^2 c^4 B}}{8m^2 c^4}$$

$$= \frac{4Ec^4(M^2 - 2m^2) \pm \sqrt{16E^2 c^8 (M^2 - 2m^2)^2 - 16m^2 c^8 (M^4 c^4 + 4m^2 (E^2 - M^2 c^4))}}{8m^2 c^4}$$

$$= E \frac{M^2 - 2m^2}{2m^2} \pm \frac{1}{2m^2} \sqrt{E^2 (M^4 - 4M^2 m^2) - m^2 M^2 c^4 (M^2 - 4m^2)}$$

$$= E \frac{M^2 - 2m^2}{2m^2} \pm \frac{M}{2m^2} \sqrt{(M^2 - 4m^2)(E^2 - m^2 c^4)}$$

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We have: $M = 10.58 \frac{\text{GeV}}{c^2}$, $m = 0.511 \frac{\text{MeV}}{c^2}$ (from Griffiths)

$$\text{So } \frac{M^2 - 2m^2}{2m^2} = \frac{1}{2} \left(\frac{M}{m} \right)^2 - 1 = \frac{1}{2} \left(\frac{10.58 \times 10^9}{0.511 \times 10^6} \right)^2 - 1 = 2.14 \times 10^8$$

$$\frac{M}{2m^2} = \frac{10.58 \times 10^9}{2 \times (0.511 \times 10^6)^2} \left(\frac{\text{eV}}{c^2} \right)^{-1} = 2.03 \times 10^{-2} \frac{c^2}{\text{eV}}$$

$$M^2 - 4m^2 \approx M^2 = 1.12 \times 10^{20} \frac{\text{eV}^2}{c^4}$$

For $E = 9.0 \text{ GeV}$, we get:

$$\bullet E \frac{M^2 - 2m^2}{2m^2} = 9.0 \text{ GeV} \times 2.14 \times 10^8 = 1.93 \times 10^{18} \text{ eV}$$

$$\bullet \frac{M}{2m^2} \sqrt{(M^2 - 4m^2)(E^2 - m^2 c^4)} \approx \frac{M}{2m^2} \sqrt{M^2 E^2} = \frac{M^2}{2m^2} E \approx \frac{M^2 - 2m^2}{2m^2} E = 1.93 \times 10^{18} \text{ eV}$$

The precision of our calculations is not sufficient to account for the difference between the two quantities, but it turns out that the solution we ~~wanted~~^{need} is actually the difference, not the sum. To see that that's true,

$$\text{first consider } E'_+ = 2 \times 1.93 \times 10^{18} \text{ eV} \approx \frac{M^2}{m^2} E$$

It's easy to see that it's not a solution of (A), though it is a solution of the equation that follows (A) — we introduced an extra solution when we squared (A).

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Indeed, the RHS (right-hand side) of (*) is

$$\text{RHS} \approx M^2 c^4 - 2E \frac{M^2}{m^2} E = \frac{M^2}{m^2} (m^2 c^4 - 2E^2) < 0$$

whereas, of course,

$$\text{LHS} > 0$$

So ~~both~~ LHS and RHS have different signs, but, when squared, produce the same result - that's how the extra solution appeared.

So what we need is E_-' . We could use a computer to calculate both quantities to a great precision and subtract them or we can be more clever and manage using only our calculator. We know that

$$m^2 \ll M^2, \quad m^2 c^4 \ll E^2,$$

so we can expand the square root:

$$\begin{aligned} \sqrt{(M^2 - 4m^2)(E^2 - m^2 c^4)} &= \sqrt{\overset{\uparrow \text{big}}{M^2 E^2} - \overset{\uparrow \text{smaller}}{(4m^2 E^2 + m^2 c^4 M^2)} + \overset{\uparrow \text{very small}}{4m^4 c^4}} \\ &= ME \sqrt{1 + \frac{-\overset{\uparrow \text{smaller}}{(4m^2 E^2 + m^2 c^4 M^2)} + \overset{\uparrow \text{very small}}{4m^4 c^4}}{M^2 E^2}} \\ &= ME \left(1 - \frac{1}{2} \frac{+4m^2 E^2 + m^2 c^4 M^2}{M^2 E^2} + \text{smaller terms} \right) \end{aligned}$$

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Then

$$E'_- = E \frac{M^2 - 2m^2}{2m^2} - \frac{M}{2m^2} ME \left(1 - \frac{1}{2} \frac{4m^2 E^2 + m^2 c^4 M^2}{M^2 E^2} \right)$$

$$= E \frac{M^2}{2m^2} - E - \frac{M^2}{2m^2} E + E + \frac{M^2 c^4}{4E} = \frac{M^2 c^4}{4E}$$

So for $E = 9.0 \text{ GeV}$, we get:

$$E' = \frac{(Mc^2)^2}{4E} = \frac{(10.58 \times 10^9 \text{ eV})^2}{4 \times 9.0 \times 10^9 \text{ eV}} = \boxed{3.11 \times 10^9 \text{ eV}} \quad \text{in SLAC}$$

The exact same story repeats for $E = 8.0 \text{ GeV}$, and so:

$$E' = \frac{(10.58 \times 10^9 \text{ eV})^2}{4 \times 8.0 \times 10^9 \text{ eV}} = \boxed{3.50 \times 10^9 \text{ eV}} \quad \text{in KEK}$$

Note that our calculations would be greatly simplified if we noticed right away that $E \gg mc^2$ and $E' \gg mc^2$, so (3) and (4) give $E \approx pc$, $E' \approx p'c$. Plugging that into (2), we'd get $cp = E - E'$. Then, using (1),

$$M^2 c^4 = E^2 - c^2 p^2 = (E + E')^2 - (E - E')^2 = 4EE'$$

and

$$\boxed{E' = \frac{M^2 c^4}{4E}}$$

The lesson is: the art of approximation is a very useful thing in physics.

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- (2) Denote: E_1, E_2, p_1, p_2 - the energies and momenta of the B-mesons,
 m_B - their mass

In the rest frame of $\Upsilon(4S)$, the conservation laws give:

$$\begin{cases} Mc^2 = E_1 + E_2 \\ 0 = p_1 - p_2 \end{cases}$$

$$p_1 = p_2 \Rightarrow E_1 = E_2 \Rightarrow E_1 = E_2 = \frac{Mc^2}{2} = \frac{10.58 \text{ GeV}}{2}$$

$$E_1 = E_2 = 5.29 \text{ GeV}$$

$$p_1 = p_2 = \frac{1}{c} \sqrt{E_1^2 - m_B^2 c^4} = \frac{1}{c} \sqrt{(5.29 \times 10^9 \text{ eV})^2 - (5279 \times 10^6 \text{ eV})^2}$$

$$= \frac{3.41 \times 10^8 \text{ eV}}{c}$$

$$p_1 = p_2 = 0.34 \text{ GeV}/c$$

- (3) If a meson lives a time τ in its rest frame it will live $\gamma\tau$ in the ~~lab~~ frame (usual time dilation). We also know that $p = \gamma m v$, $E = \gamma m c^2$, so it will travel
- $$d = \gamma\tau v = \tau \gamma v = \tau \frac{p}{m_B} = 1.54 \times 10^{-12} \text{ s} \frac{0.34 \times 10^9 \text{ eV}}{3 \times 10^8 \text{ m/s}} \frac{1}{5279 \times 10^6 \text{ eV}} c^2$$

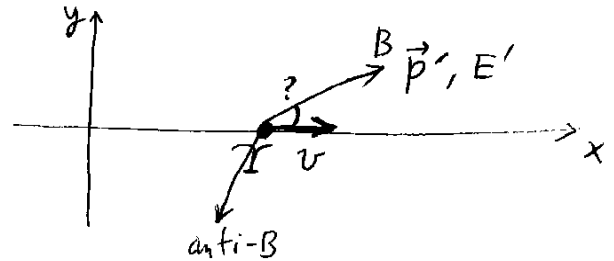
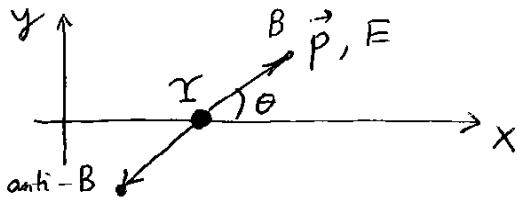
$$d = 2.97 \times 10^{-5} \text{ m}$$

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(4) Denote: γ - the dilation factor for the moving $\mathcal{I}(45)$

K - rest frame of $\mathcal{I}(45)$:

K' - lab frame:



p - the momentum of B in the $\mathcal{I}(45)$ frame
 p' - the momentum of B in the lab frame, etc.

So in the K -frame,

$$\vec{p} = (p \cos \theta, p \sin \theta, 0)$$

and the energy-momentum four-vector for B is $\begin{pmatrix} E/c \\ p \cos \theta \\ p \sin \theta \\ 0 \end{pmatrix}$

Therefore in the K' -frame the energy-momentum four-vector for B will be (by the usual transformation law)

$$\begin{pmatrix} E'/c \\ p'_x \\ p'_y \\ p'_z \end{pmatrix} = \begin{pmatrix} \gamma + \gamma\beta & 0 & 0 & 0 \\ +\gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} E/c \\ p \cos \theta \\ p \sin \theta \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{\gamma E}{c} + \gamma\beta p \cos \theta \\ +\gamma\beta \frac{E}{c} + \gamma p \cos \theta \\ p \sin \theta \\ 0 \end{pmatrix}$$

+ , not - , because K' is moving to the left w/r to K

(5) In problem 3, we found an easy formula for the distance traveled. In the lab frame, a meson will live $\tilde{\gamma} \tau$, so it will travel a distance $d = \tilde{\gamma} \tilde{v} \tau = \tilde{\gamma} m_B \tilde{v} \frac{\tau}{m_B} = \tilde{p} \frac{\tau}{m_B}$, where \tilde{p} - the momentum of the meson in the lab frame.

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In the previous problem, we found the expression for the four-momentum of the B-meson as a function of θ , so

$$\tilde{P} = \sqrt{\left(+\gamma\beta\frac{E}{c} + \gamma p \cos\theta\right)^2 + p^2 \sin^2\theta},$$

where p, E - the momentum and energy of the B-meson in the frame of $\mathcal{I}(4S)$ was calculated in problem 2 and γ, β - the dilation factors for the moving $\mathcal{I}(4S)$ can be obtained from the results of problem 1 after a few additional calculations.

We have for the average momentum:

$$\begin{aligned} \langle \tilde{P} \rangle &= \int_{\cos\theta=-1}^{\cos\theta=1} \tilde{P}(\theta) dP(\theta) = \int_{-1}^1 \sqrt{\left(+\gamma\beta\frac{E}{c} + \gamma p \cos\theta\right)^2 + p^2(1-\cos^2\theta)} \frac{3}{8}(1+\cos^2\theta) d\cos\theta \\ &= \frac{3}{8} \int_{-1}^1 \sqrt{\left(+\gamma\beta\frac{E}{c} + \gamma p x\right)^2 + p^2(1-x^2)} (1+x^2) dx \end{aligned}$$

It is possible to take this integral analytically (integrals of the form $\int R(\sqrt{ax^2+bx+c}, x) dx$, where R is any rational function of two variables, can be taken, for example, by completing the square, getting it into one of the ~~two~~ three forms: $R(x, \sqrt{a^2-x^2})$, $R(x, \sqrt{x^2-a^2})$, $R(x, \sqrt{x^2+a^2})$, using the substitutions $x = a \sin t$, $x = a \overset{\cosh}{\cosh} t$, $x = a \sinh t$ respectively, then using the substitutions $y = \tan \frac{t}{2}$, $y = \tanh \frac{t}{2}$, $y = \tanh \frac{t}{2}$ respectively),

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but it is more convenient to use a graphing calculator or a computer program. In problem 2, we found that

$$p = 0.34 \frac{\text{GeV}}{c}, \quad E = 5.29 \text{ GeV}$$

In problem 1, we found the energy of the positron for the two values of the energy of the electron. The total energy of $\Upsilon(4S)$ is the sum of the energies of the electron and positron, so

$$\text{for SLAC, } E = 9.0 \text{ GeV} + 3.11 \text{ GeV} = 12.11 \text{ GeV}$$

$$\gamma = \frac{E}{Mc^2} = \frac{12.11 \text{ GeV}}{10.58 \text{ GeV}} = 1.14$$

$$\gamma = \frac{1}{\sqrt{1-\beta^2}}, \quad \beta = \sqrt{1-\frac{1}{\gamma^2}} = 0.49$$

$$\gamma\beta \frac{E}{c} = 2.95 \frac{\text{GeV}}{c},$$

$$\gamma p = 0.39 \frac{\text{GeV}}{c}$$

$$p^2 = 0.12 \left(\frac{\text{GeV}}{c}\right)^2$$

$$\text{for KEK, } E = 8.0 \text{ GeV} + 3.5 \text{ GeV} = 11.5 \text{ GeV}, \quad \gamma = \frac{11.5}{10.58} = 1.09$$

$$\gamma\beta = \sqrt{\gamma^2 - 1} = 4.26, \quad \gamma\beta \frac{E}{c} = 2.25 \frac{\text{GeV}}{c}$$

$$\gamma p = 0.37 \frac{\text{GeV}}{c}, \quad p^2 = 0.12 \left(\frac{\text{GeV}}{c}\right)^2$$

We plug these values into Mathematica to find:

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for SLAC, $\langle \tilde{p} \rangle = 2.96 \frac{\text{GeV}}{c}$

for KEK, $\langle \tilde{p} \rangle = 2.27 \frac{\text{GeV}}{c}$

Then the average distance the meson will travel in the

lab is

$$\langle d \rangle = \langle \tilde{p} \rangle \frac{\tau}{m_B} = \begin{cases} 2.96 \frac{\text{GeV}}{c} \frac{1.54 \times 10^{-12} \text{s}}{5279 \text{ MeV}/c^2} = \boxed{2.59 \times 10^{-4} \text{ m, SLAC}} \\ 2.27 \frac{\text{GeV}}{c} \frac{1.54 \times 10^{-12} \text{s}}{5279 \text{ MeV}/c^2} = \boxed{1.99 \times 10^{-4} \text{ m, KEK}} \end{cases}$$